THE ASYMPTOTICS AND THE DESCENDANTS OF THE 3D-INDEX

RENEW QUANTUM SEMINAR

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References:

2209.02843 Periods, the mero 3D-index and TV-invariants 2301.00098 The descendants of the 3D-index

## Contents:

- The 3d-index and the state-integral • The 3d-index (definitions, properties)
- · Descendants (algebra)
- · Asymptotics (analysis)

The 3d-index and the state integral 3d-3d correspondence on one side gives an N=2 superconformal field theory.  $T_{M}[S^{3}]$  on  $S^{3}=D^{2}xS'\cup D^{2}xS'$   $S=\binom{D-1}{1}eSLZ$ or on  $T_{M}[s^{2}xs']$  on  $s^{2}xs' = D^{2}xs' \cup D^{2}xs' \in [(::,)eal_{2}Z]$ and more generally on ye SL2Z TM[Lens space] Λ

• 
$$T_{M}[S^{3}] = Andersen-Kashaev stateintegral, expressed(conjecturely)bilinearly in terms of functions of q and  $\tilde{q}$   
where  $q = e^{2\pi i \tau}$ ,  $\tilde{q} = e^{-2\pi i / \tau}$$$

• 
$$T_{M}$$
 [lens space] is a state-integral  
defined using an  $SL_{Z}$  extension of  
the Faddeev quantum dilogarithm [GKZ]  
expressed bilinearly in terms of the  
same functions of q and  $q^{H}$   
 $q=e^{2\pi i \tau}$ ,  $q^{H}=e^{2\pi i g(\tau)}$ ,

## The DGG 3d-index

Building block: fetrahedron index  

$$I_{\Delta}(m,e)(q) = \sum_{n=te}^{\infty} (-1)^{n} \frac{q^{\frac{1}{2}n(n+1)-(n+\frac{1}{2}e)m}}{(q;q)n(q;q)_{n+e}} \in \mathbb{Z}(q)|^{n}$$

$$(m)_{t} = \max\{0, m\}. \quad The partition function of I_{\Delta}^{rot}[s^{2}xs']$$

$$deg(I_{\Delta}(m,e)(q) = \frac{1}{2}(m_{t}(m+e)_{t}+(-m)_{t}e_{t}+(-e)_{t}(-e-m)_{t}+\max\{0,m,-e]) \ge 0$$



Eliminate z' using 
$$z'=-\frac{1}{zz''}$$
  

$$\sum_{j=1}^{N} A_{ij} \log_{2} + B_{ij} \log_{2} j''= \pi i B_{i} \quad i=1...,N+2$$
Let  $A = (a_{1}|...|a_{N}) \quad B_{\pi}(b_{1}|...|b_{N})$   
So  $T \longrightarrow (A|B|U)$   
Def  $I_{\tau}^{rot}[s' \times s^{2}] \quad computed via 3J-3d \quad correspondence, a$   
Def  $I_{\tau}^{rot}(n,n')(q) = \sum_{T} (k,n,n')(q)$   
 $k \in \mathbb{Z}^{N}$   
 $S(k,n,n')(q) = (-q'/L)^{V,k-(n-n')V_{\lambda}} g_{N}(n+n')V_{\lambda}$   
 $\cdots \prod_{j=1}^{N} I_{\lambda}(\lambda_{j}^{"}(n-n')-b_{j}\cdot k, -\lambda_{j}\cdot (n-n')+a_{j}\cdot k)(q)$   
 $Thm [GHRS](Allf T=1-elhicient (ie has no nonperipheral Normal tori, and no normal  $s^{-1}$ ) then  
 $I_{\tau}^{rot}(n,n')(q) \in \mathbb{Z}(l(q))$  is well-defined  
(b) It is a topological invariant of cusped hyperbolic manifolds  
 $Thm [GK]$  Alternative proof of topological invariance Using state-integral formula  $I_{\tau}^{rot}$$ 

In pathicular, 
$$I_{T}^{mer}(0,0)[q] = \sum_{r \in \mathbb{Z}} I_{T}^{rot}(n,n)[q]$$
  
Properties of 3d-index (Conjectural)  
Factorization  
 $I_{T}^{rot}[s' \times s^{2}] = \langle I_{T}^{rot}[s' \times D^{2}], I_{T}^{rot}[s' \times D^{2}] \rangle$   
 $I_{T}^{rot}[q] = H_{T}(q)B_{T}H_{T}(q^{-1})^{T}$   
 $\mathbb{Z} \times \mathbb{Z}$   $\mathbb{Z} \times r \times r \times r \times \mathbb{Z}$   
 $H_{T}(q) = (h_{T}^{(\alpha)}(q))$   $\alpha = 1 - r, n \in \mathbb{Z}$   
 $\int_{T} \log e H(q)$  is a properly normalized  
hundamental solution of liveor q-difference  
 $equation$ 

$$\frac{\operatorname{Regularity}}{\operatorname{I}_{\tau}^{rot}(n,n')(q)} = \lim_{x \to 1} \sum_{\alpha} B_{\tau}^{(\alpha')}(q^{n'}x^{-'};q^{-'}) B_{\tau}^{(\alpha')}(q^{n}x;q)$$
where  $B_{\tau}^{(\alpha)}(x;q)$  are x-deformed holo blocks

Descendants  
Insertions, defects, line operators  

$$M = S^{2} - K$$
  $\longrightarrow$   $O \in M(T)$   
LCM  
 $M(T) = \frac{W_{q}(T)}{(W_{q}(T) \cdot \text{Lagrangians} + \text{edge eqns} \cdot W_{q}(T))}$   
 $W_{q}(T) = \mathcal{Q}(q) < \hat{z}_{j}, \hat{z}'_{j}, \hat{z}'_{j}, \hat{z}''_{j} | i=1...N > = q - Weyl algebra.$   
 $\hat{z} \hat{z}' = q \hat{z}' \hat{z} + \text{cyclic permutations}$   
 $\hat{z} \hat{z}' \hat{z}'' = -q$ 

$$\begin{split} \underbrace{\operatorname{Vel}}_{t,0} & \text{ If } 0 = \prod_{j=1}^{n} \hat{z}_{j}^{\alpha_{j}} \hat{z}_{j}^{\gamma_{j}} \hat{b}_{j} \\ & \operatorname{I}_{t,0}^{rot}(n,n!(q)) = \sum_{k \in \mathbb{Z}^{N}} (0 \circ S_{t})(k,n,n!(q)) \\ & \operatorname{Vel}_{t,0}(n,n!(q)) = (-q^{1/2})^{\nu_{k}} \cdot (n-n!)^{\nu_{k}} q^{k_{k}}(n+n!)^{k_{k}} + L_{0}(n,n!(k)) \\ & \cdot \prod_{j=1}^{n} I_{k}(A_{j}^{\gamma_{k}}(n-n!) - b_{j} \cdot k \cdot b_{j}; -\lambda_{j}(n-n!) + a_{j} \cdot k - \alpha_{j})(q) \\ & \underbrace{\operatorname{Node}}_{j=1} The inservion \mathcal{E}_{i} of the ith edge acts as \\ & (\mathcal{E}_{i} \circ S_{T})(k,n,n') = qS_{T}(k-e_{i},n,n') \\ & \operatorname{hence}_{inviniting} \text{ over } ke\mathbb{Z}^{N}, \text{ if follows } \mathcal{E}_{i}:-q \\ & annihilates I^{ret}(n,n')(q). \\ & \underbrace{\operatorname{Properties}_{i \in S} of descendent 3d-index(Conjecture)}_{T,0}(a) = I_{T,0}^{rot}(a) = H_{T,0}(a)B_{T}H_{T}(a^{-1})^{t} \\ & (b) There exists Q_{T,0}(a) = M_{T}(a^{-1})^{t} \\ & (b) There exists Q_{T,0}(a) = M_{T}(a^{-1})^{t} \\ & I_{T,0}^{rot}(r^{-1}) = Q_{T,0}(a) I_{T}^{ret}[r^{-1}] \\ & \operatorname{Cor}_{T,0}(a) = i \text{ uniquely determined by} \\ \end{aligned}$$

.

(a) rxr matrices 
$$I_{\tau}^{rot}(q)[r]$$
 and  $Q_{\tau}(q)$   
and  $(vacum)$   $\tau,0$   
(b) peir of linear q-difference eqns  
 $\hat{A}_{\tau}$  and  $\hat{A}_{\tau,0}$ .  
Def let  $DI_{\tau,0}^{rot} = \text{Span} \left\{ I_{\tau,0}^{rot}(n,n')(q) \right\} n,n' \in \mathbb{Z} \right\}$   
 $(a \text{ hin. dimensional } O(q^{v_1}) - vct \text{ space})$   
 $(a \text{ hin. dimensional } O(q^{v_1}) - vct \text{ space})$   
 $Cor$   $U$   $DI_{\tau,0}^{rot} = DI_{\tau}^{rot}$   
 $O \in \mathcal{U}(\tau)$   $T, 0$   $T$ .  
In other words, the descendants of the rotated  
 $T + index$  are expressed effectively by a finite  
size matrix over  $O(q^{v_1})$ .  
 $Conjective = \hat{A}_{\tau,0}(M,L,q) = \hat{A}_{\tau}(M,L)$ 

Examples  
4. knot  

$$G_{1} = \begin{pmatrix} 2 & 2 \\ 0 & 0 \\ 1 & 1 \end{pmatrix} G_{1}' = \begin{pmatrix} 1 & 1 \\ 1 & 1 \\ 0 & 0 \\ 1 & 1 \end{pmatrix} G_{1}' = \begin{pmatrix} 0 & 0 \\ 2 & 2 \\ 0 & -1 \\ 1 & -3 \end{pmatrix}$$

$$A = \begin{pmatrix} 1 & 1 \\ 1 & 0 \end{pmatrix} B = \begin{pmatrix} -1 & -1 \\ 0 & -1 \end{pmatrix} V = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

$$I_{U_{1}}^{vot}(n, n')(q) = \sum_{k_{A}} q^{k_{A}} \begin{pmatrix} n+n' \\ 2 \end{pmatrix} I_{k_{A}}(k_{A}, k_{A} + k_{A})(q) I_{A}(k_{A} - n+n', k_{A} - n+n')(q)$$

$$k_{A}, k_{A} \in \mathbb{Z}$$

$$\frac{Factorization}{I^{rot}(n,n')(q) = -\frac{1}{2}h_{n'}^{(1)}(q^{-1})h_{n'}^{(0)}(q) + \frac{1}{2}h_{n'}^{(0)}(q^{-1})h_{n'}^{(1)}(q)}{h_{n'}(q)}$$

$$h_{n}^{(0)}(q) = (-1)^{n}q^{\ln|(2|n|+1)/2} \sum_{k=0}^{2} (-1)^{k} \frac{q^{k(k+1)/2} + \ln|k}{(q;q)_{k}(q;q)_{k+2|n|}}$$

$$h_{n}^{(1)}(q) = sim:|ar, a bit$$

$$more \ Gunplicated$$

$$satisfy \ linear \ q-dilference \ eqn \ (neZ, \alpha=0,1):$$

$$q^{2+2n}(q^{3+2n}-1)h_{n'}^{(\alpha)}(q) + ... + q^{3+2n}(q^{1+2n}-1)h_{n+2}^{(\alpha)}(q) = 0$$

$$\int \alpha ti f y$$

$$I_{\hat{Q}}^{rot}(n,n')(q)[2] = \frac{1}{1-q} \begin{pmatrix} q-2 & q^{\frac{1}{2}} \\ -q^{\frac{1}{2}} & q+1-q^{\frac{1}{2}} \end{pmatrix} \quad I^{rot}(n,n')(q)[2] + O(q^{121})$$
  
illerstrating the dramatic cancellation of the  
Sum of products of q-series into a short rat hunchon.

In particular  

$$I_{\hat{0}}^{rot}(0,0)(q) = \frac{1}{1-q} ((q-2) I^{rot}(0,0)(q) + q^{\frac{1}{2}} I^{rot}(0,1)(q)) . \square$$

This was repeated for 
$$5_2$$
 knot  $(r=3, 3 \text{ tetrahedra})$   
and even more impressively for the  $(-2,3,7)$   
preted knot  $(r=6, 3 \text{ tetrahedra})$ 

where I<sup>rot</sup>(q)[6] has 20 digits integer Coefficients for the coeff of q 160, and so did I<sup>rot</sup>(q)[6], however their ratio is a short rational function with coefficients integers between -5,5.

## Asymptotics

The story begins in 2011 on a train to Bonn with Don Zagier after a Gnference in the Diablerets

$$g(q) = \sum_{n=0}^{\infty} (-1)^n \frac{q^n \frac{n(n+1)}{2}}{(q;q)_n^2} = 1 - q - 2q^2 - 2q^3 - 2q^4 + q^6 + \dots$$

$$(q;q)_n = (1 - q)(1 - q^2) \dots (1 - q^n)$$

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$$g(e^{\frac{2\pi}{N}}) \sim \frac{1}{\sqrt{N}} \frac{\sqrt{2}}{\sqrt{3}} \left( \cos \frac{v_{ol}}{2\pi} N + \sin \frac{v_{ol}}{2\pi} N \right) \left( 1 - \frac{11}{72\sqrt{3}} \frac{2\pi}{N} + \frac{697}{2(7273)^4} \left( \frac{2\pi}{N} \right)^{4} \right)$$
  
First the period  $\frac{v_{ol}}{2\pi} \sim 0.32$  of the oscillation  
was guessed, then 697 was found by search in  
files regarding the volume Gnjecture of the 4, knot.  
 $v_{ol} = v_{ol}(4_1) = 2 \operatorname{Imbig}(e(\frac{1}{6}))$ 

Thus  

$$g(ett) \sim J_{e} \sqrt{\tau} \left( \hat{\Phi}(2n;\tau) - i \hat{\Phi}(-2n;\tau) \right)$$
where  $\hat{\Phi}(h) = e^{\frac{i \sqrt{n} (l+1)}{h}} O(h)$   

$$\Phi(h) = \frac{1}{\sqrt{\tau_{s}}} \left( 1 - \frac{1!}{72\sqrt{\tau_{3}}} h + \frac{697}{2(72\sqrt{\tau_{3}})^{s}} h^{2} + \frac{724351}{30(72\sqrt{\tau_{3}})^{s}} h^{s} + \dots \right)$$
satisfies  

$$< 4_{1} \sum_{N} \sim N^{3/2} \hat{\Phi}\left(\frac{2ni}{N}\right)$$

$$< 4_{1} \sum_{N} = Kash aev invariant of 4_{1}$$

$$\cdot \tau \rightarrow 0 \quad \text{horizontally ocarg } \tau < \epsilon \quad \text{fixed}$$
Here  $\hat{\Phi}(2n;\tau)$  is exponentially bigger than  $\hat{\Phi}(-2n;\tau)$   
but now  $\tilde{q}$  corrections are visible,  $\tilde{q} = e^{-2ni/\tau}$   
and in fact  

$$g(e(\tau)) \sim J_{s} \sqrt{\tau} g(e(-\frac{1}{\tau})) \hat{\Phi}(2n;\tau)$$
Actually this needs to be interpreted  
eq via median Borel sourmation. In this  
Case  $\hat{\Phi}(-2n;\tau) \ k(e(-\sqrt{\tau})) \ also \ contributes \ and$   
then have identity.  
Conclusion Radial asymptotics of  $q$ -hypergeometric  
Series at  $q$ -wood of unity depends on sectors.

Fast forward to last year.

 $I_{\tau}^{\text{mer}}(0,0)(q) = \sum_{\tau} I_{\tau}^{\text{ret}}(n,n)(q)$ (the relation between the [GK] 3d-index and the rotated 3d-index?  $I_{4_{1}}^{mer}(0,0)(q) = 1 - 4q - q^{2} + 36q^{3} + 70q^{4} + 100q^{5} + 34q^{6} - 116q^{7} + \cdots$ when q=e(t), t ->0 nearly horizontally we found that  $I_{u_1}^{mer}(o, o)(q) \sim e^{\frac{2 vol}{2 n \tau}} \frac{1}{2^{3} 4 2^{1/2} \sqrt{\tau}}$  $\times \left( 1 - \frac{19}{94\sqrt{23}} 2\pi i \tau + \frac{1333}{1152\sqrt{23}} (2\pi i \tau)^2 - \frac{1601717}{11020\sqrt{23}} (2\pi i \tau)^3 + \frac{1333}{1152\sqrt{23}} \right)$ Then, we recognized the number 1333 appearing in the asymptotics of the TV invariant TV4, m+ replacing m+ 2 by T. Then we checked 7 more coefficients.  $\frac{\text{Def } TV}{K, m+\frac{1}{2}} = * \sum_{k=1}^{\infty} \left[ k J^2 \left| J_{K, k}(q) \right|^2 \right|_{q=e(\frac{1}{m+1})}$ Chen-Yang conjectived that TV grows exponentially Capturing the volume of a hyperbolic knot.

$$\frac{Conjecture}{k} [GW] \quad \tau \rightarrow 0 \text{ rearly harizontally} \\ \frac{T^{mer}_{k}(0,0)(e(\tau))}{k} \approx \sum_{k \in \mathbb{Z}} \widehat{\Phi}_{k,n}^{(\sigma_{1})}(2ni\tau) \widehat{\Phi}_{k,n}^{(\sigma_{n})}(-2\pii\tau) \\ \frac{TV_{k,n+\frac{1}{2}}}{k} \approx \sum_{k \in \mathbb{Z}} \widehat{\Phi}_{k,n}^{(\sigma_{1})}\left(\frac{2ni}{m+\frac{1}{2}}\right) \widehat{\Phi}_{k,n}^{(\sigma_{n})}\left(-\frac{2\pii}{m+\frac{1}{2}}\right) \\ \frac{T^{rot}_{k}(n,n')(e(\tau))}{k} \approx \widehat{\Phi}_{k,n}^{(\sigma_{1})}(2ni\tau) \widehat{\Phi}_{k,n}^{(\sigma_{n})}(-2\pii\tau) \\ \frac{T^{rot}_{k}(n,n')(e(\tau))}{k} \approx \widehat{\Phi}_{k,n}^{(\sigma_{1})}(2ni\tau) \widehat{\Phi}_{k,n'}^{(\sigma_{n})}(-2\pii\tau)$$

where 
$$\hat{\Phi}_{n}^{(0)}(h) = e^{V_{T}} \Phi_{n}^{(r)}(h)$$
  
 $\Phi_{n}^{(0)}(h) = \frac{1}{V_{ST}} (1 + F_{T}[n,h])$   
 $\sigma = \sigma_{1} = \text{geometric representation}$   
 $\sigma = \sigma_{2} = \text{Guplex Gujugate}$   
and  $(\hat{\Phi}_{n}^{(0)}(h))$  is a fundamental solution  
to the  $\hat{A}$ -eqn annihilating  $I^{rot}(q)$ .  
 $\hat{\Phi}_{n}^{(0)}(h)$  can be computed (once  $\hat{A}$  is known)  
by WKB.  
This Gnjechred was Gnhimed for 4, and 52  
in [GW].

Vertical asymptotics  
Conjecture [GW] When two vertically  
Irot (n,n') (e(t)) ~ 
$$\sum \varepsilon_0 \hat{\Phi}_{n}^{(\sigma)}(2nit) \hat{\Phi}_{n'}^{(G)}(-2nit)$$
  
where  $\varepsilon_0 = \begin{bmatrix} -1 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 \end{bmatrix}$   
Note that after cancelling the exponential  
lactors, the Gustaut terms of the power series  
are possitive numbers. The asymptotic series  
are  $0(-t^{poly})$  perhaps a ensequence of  
Unitarity.  
Now vertical asymptotics  $I_{41}^{mer}(0,0)(q)$   
Numerically, we found out that  
 $I_{41}^{mer}(0,0)(q) \sim e^{\frac{2\pi i}{2\pi t}} \frac{1}{152t-3}e^{(2nit)^2} - \frac{160(7i)^2}{414780t-3}(2nit)^2 + \dots \right)$   
 $-i \times (replace t by  $-t$  in  $J$ )  
 $t \chi_{ij} - \frac{i}{t} + \chi'_{ij}$  it  $t + \dots$   
where  $\chi_{ij} = 0.44582579449a35614a77$   
From previous experience we tried to reagginze  
this or an alg number and n but failed.  
Then Compbell suggested to look for periods$ 

of A-poly curve  

$$y = x^{-2} x^{-1} - 2x + x^{2} \quad (\text{elliptic curve})$$
and much to our surprise  

$$x_{u_{1}} = \frac{\partial u_{1} + \partial u_{1}}{2\pi}$$

$$\partial_{u_{1}} = \int \frac{du}{\sqrt{e^{-2u} - 2e^{-u} - 1 - 2e^{u} + e^{2u}}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{u} + e^{2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{u} + e^{2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{u} + e^{2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{u} + e^{2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{u} + e^{2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{u} + 2e^{-2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{u} + 2e^{-2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{u} + 2e^{-2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{u} + 2e^{-2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{u} + 2e^{-2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{u} + 2e^{-2u}} \cdot \frac{1}{2e^{-2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{u} + 2e^{-2u}} \cdot \frac{1}{2e^{-2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{u} + 2e^{-2u}} \cdot \frac{1}{2e^{-2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{u} + 2e^{-2u}} \cdot \frac{1}{2e^{-2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{-u} + 2e^{-2u}} \cdot \frac{1}{2e^{-2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{-u} + 2e^{-2u}} \cdot \frac{1}{2e^{-2u}} \cdot \frac{1}{2e^{-2u} - 2e^{-u} - 1 - 2e^{-u} + 2e^{-2u}} \cdot \frac{1}{2e^{-2u}} \cdot \frac{1}{2e^{-2u$$

Periods are Mellin-Burnes integrals  
Thm [GK]  

$$\frac{1}{2\pi i} \int_{\xi-iR} \frac{B(z,z)^2}{(\cos\pi z)^2} dz = 2 \int_{-\infty}^{1} \frac{dx}{\sqrt{(l-x)(l-x+4x^2)}} = 5.60241216$$
where  

$$B(x_1y) = \frac{\Gamma(x)\Gamma(y)}{\Gamma(x+y)} = \int_{0}^{1} t^{x} (l-t)^{y} \frac{dt}{t(l-t)}$$
if the Eyler B-Kinchion.